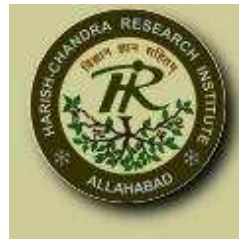

Signatures of non-universal Gaugino and Scalar masses at the Large Hadron Collider (LHC)

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Work done with Aresh Krishna Datta, and Biswarup Mukhopadhyaya

(arXiv:0708.2427, JHEP 10(2007)080)

(arXiv:0804.4051)

Plan of the talk

- Basic ideas behind the work.
- Non-universal gaugino masses
 - Model
 - Choice of Parameters
 - Collider simulation
 - Results
- Non-universal Scalar masses
 - Squark-Slepton non-universality
 - Third family scalar non-universality
 - Scalar non-universality for $SO(10)$ $D - term$

Basic ideas behind the work : Gaugino mass non-universality

- Theoretical prediction \longrightarrow SUSY GUT \longrightarrow Non-universal gaugino masses at M_{GUT} .
- Non-universality of gaugino masses \longrightarrow affects the chargino-neutralino mass composition.
- Collider signature \longrightarrow should get altered for various gaugino non-universal ratios.
- Intention \longrightarrow distinguish them at LHC.
- To do so \longrightarrow perform collider simulation for 'multichannel search' over a wide region of *parameter space*

Basic ideas behind the work: Scalar mass non-universality

- Theoretical SUSY GUT & Various phenomenological models → Non-universal scalar masses at M_{GUT}
- Non-universal scalar masses → alters scalar mass hierarchy → alters decay branching fractions.
- Intention → distinguish scalar mass scenarios at LHC.
- To do so → perform collider simulation for 'multichannel search' over a wide region of *parameter space*

Non-universal gaugino mass: Model

- Theoretical framework: $N=1$ Supergravity embedded in $SU(5)$ or $SO(10)$ GUT group.
- Gaugino masses depend crucially on →
 - Gauge kinetic function $f_{\alpha\beta}(\Phi)$.
 - Analytic function of chiral superfields Φ_i .
 - Transforms as → symmetric product of the adjoint representation

Non-universal gaugino mass: (Model contd.)

- Part of the **N=1 supergravity lagrangian** containing *kinetic energy and mass terms for gauginos and gauge bosons*

$$\begin{aligned} e^{-1} \mathcal{L} &= -\frac{1}{4} \text{Re} f_{\alpha\beta}(\phi) (-1/2 \bar{\lambda}^\alpha \not{D} \lambda^\beta) \\ &\quad - \frac{1}{4} \text{Re} f_{\alpha\beta}(\phi) F_{\mu\nu}^\alpha F^{\beta\mu\nu} \\ &\quad + \frac{1}{4} e^{-G/2} G^i ((G^{-1})^j_i) [\partial f_{\alpha\beta}^*(\phi^*) / \partial \phi^{*j}] \lambda^\alpha \lambda^\beta + h.c \end{aligned}$$

where

- $G^i = \partial G / \partial \phi_i$ and $(G^{-1})^j_i$ is the inverse matrix of $G^j_i \equiv \partial G / \partial \phi^{*i} \partial \phi_j$,
- λ^α is the **gaugino field**, and
- ϕ is the scalar component of the chiral superfield Φ .

Non-universal gaugino mass: (Model contd.)

- In terms of the **non-singlet** Φ^N fields :

$$f_{\alpha\beta}(\Phi^j) = f_0(\Phi^S)\delta_{\alpha\beta} + \sum_N \xi_N(\Phi^S) \frac{\Phi^N_{\alpha\beta}}{M} + \mathcal{O}\left(\frac{\Phi^N}{M}\right)^2$$

where

- f_0 and ξ^N are functions of chiral singlet superfields and
 - $M = M_{Pl}/\sqrt{8\pi}$.
- Contribution to $f_{\alpha\beta}$ from Φ^N 'has to come' through symmetric products of the adjoint representation of associated GUT group.

Non-universal gaugino mass: (Model contd.)

- For SU(5) possible non-singlet irreducible representations (to which Φ^N can belong) :

$$(24 \times 24)_{symm} = 1 + 24 + 75 + 200$$

- For SO(10) :

$$(45 \times 45)_{symm} = 1 + 54 + 210 + 770$$

Non-universal gaugino mass: (Model contd.)

- To obtain low energy effective theory \longrightarrow replace Φ^S and Φ^N by their vev's and get $\langle f_{\alpha\beta} \rangle$.
- $\langle f_{\alpha\beta} \rangle$ get the form $f_{\alpha}\delta_{\alpha\beta} \longrightarrow$ Non-Universal
- The value of $\langle f_{\alpha\beta} \rangle \longrightarrow$ crucially depends on the **specific representation responsible for the process.**
- If symmetry breaking occurs via **gauge singlet fields only**
 $\longrightarrow f_{\alpha\beta} = f_0\delta_{\alpha\beta} \longrightarrow \langle f_{\alpha\beta} \rangle = f_0 \longrightarrow$ Universal

Non-universal gaugino mass: (Model contd.)

- Simplify \longrightarrow Neglect the non-universal contributions to the gauge couplings at the GUT scale.

Table 1: Gaugino mass ratios for SU(5).

Representation	$M_3 : M_2 : M_1$ at M_{GUT}	$M_3 : M_2 : M_1$ at M_{EWSB}
1	1:1:1	6:2:1
24	2:(-3):(-1)	12:(-6):(-1)
75	1:3:(-5)	6:6:(-5)
200	1:2:10	6:4:10

Non-universal gaugino mass: (Model contd.)

- Considered only the lowest representation (54) of SO(10).

Table 2: Gaugino mass ratios for SO(10).

Representation	$M_3 : M_2 : M_1$ at M_{GUT}	$M_3 : M_2 : M_1$ at M_{EWSB}
1	1:1:1	6:2:1
54(i): $H \rightarrow SU(2) \times SO(7)$	1:(-7/3):1	7:(-5):1
54(ii): $H \rightarrow SU(4) \times SU(2) \times SU(2)$	1:(-3/2):(-1)	7:(-3):(-1)

(Chamoun et al. Nucl.Phys.B 624(2002)81)

Choice of SUSY parameters: Two options explored

- pMSSM: A phenomenological model
 - low-energy scalar masses **phenomenological**
→ **degenerate squark and slepton mass** with non-universal M_i at M_{GUT} .
- Non-universal SUGRA:
 - Generated from m_0 , A_0 and $sgn(\mu)$, with non-universal M_i at M_{GUT} .

Non-universal gaugino masses: Collider simulation

- The channels searched for: (ℓ stands for e or μ)
 - Opposite sign dilepton (OSD) :
 $(\ell^\pm \ell^\mp) + (\geq 2) \text{ jets} + E_{\cancel{T}}$
 - Same sign dilepton (SSD) : $(\ell^\pm \ell^\pm) + (\geq 2) \text{ jets} + E_{\cancel{T}}$
 - Single lepton* ($1\ell + \text{jets}$): $1\ell + (\geq 2) \text{ jets} + E_{\cancel{T}}$
 - Trilepton ($3\ell + \text{jets}$): $3\ell + (\geq 2) \text{ jets} + E_{\cancel{T}}$
 - Hadronically quiet trilepton* ((3ℓ)): $3\ell + 0 \text{ jets} + E_{\cancel{T}}$
 - Inclusive jet (jets): $(\geq 3) \text{ jets} + E_{\cancel{T}}$

Non-univ gaugino masses: (Collider simulation contd.)

- $E_{\cancel{E}} \geq 100 \text{ GeV}$.
- $p_{T\ell} \geq 20 \text{ GeV}$ and $|\eta_\ell| \leq 2.5$.
- An **isolated lepton** should have
 - lepton-lepton separation $\Delta R_{\ell\ell} \geq 0.2$
 - lepton-jet separation $\Delta R_{\ell j} \geq 0.4$
 - the energy deposit due to jet activity around a lepton E_T within $\Delta R \leq 0.2$ of the lepton axis should be $\leq 10 \text{ GeV}$.
- $E_{Tjet} \geq 100 \text{ GeV}$ and $|\eta_{jet}| \leq 2.5$.

Non-univ gaugino masses: (Collider simulation contd.)

- SM Background:
 - All dominant standard model (SM) events generated.
 - $t\bar{t}$ production most serious.
 - In the histograms where any of the entries in the ratio has $\sigma = S/\sqrt{B} \leq 2$ for $300\text{fb}^{-1} \rightarrow$ specially marked with a '#'.

Non-universal gaugino masses: Results

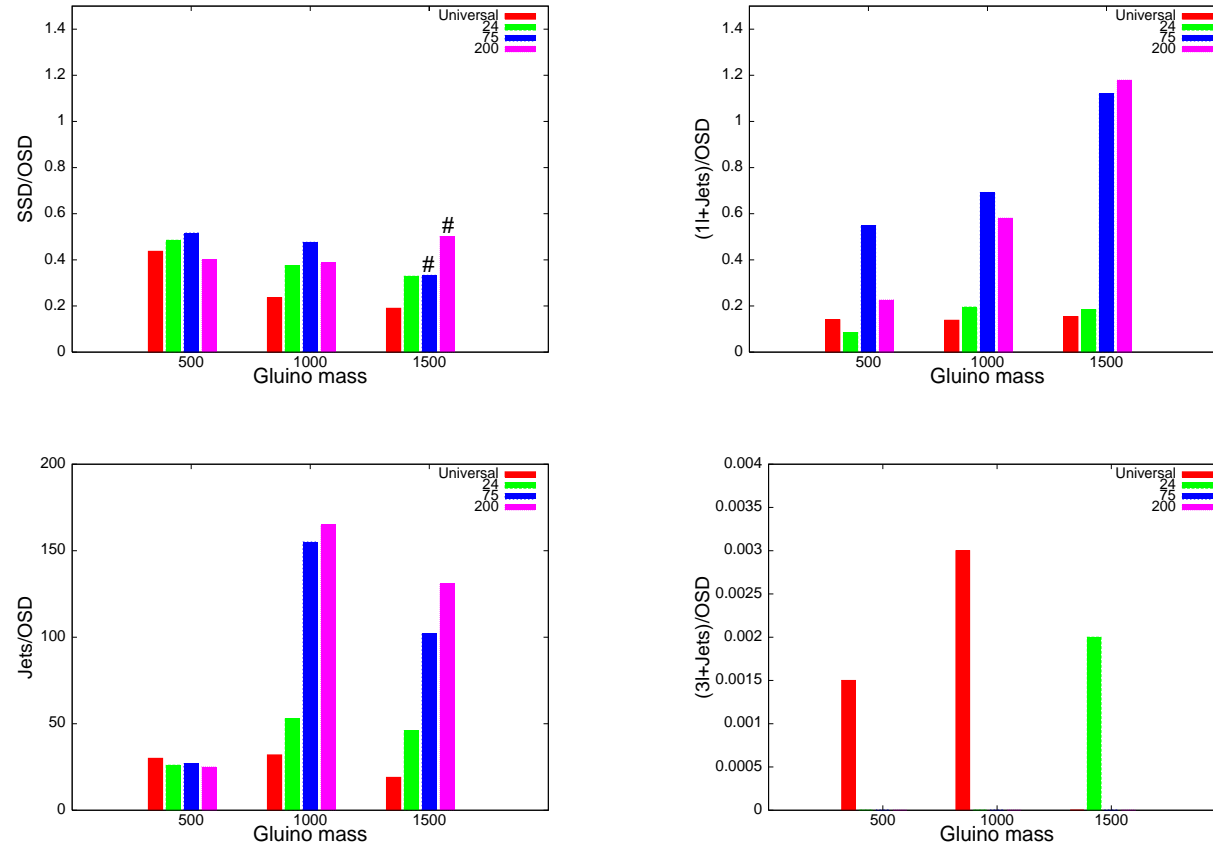


Figure 1: Event ratios for pMSSM in SU(5): $m_{\tilde{g}} = 500$ GeV, $\mu = 300$ GeV, $\tan \beta = 40$

Non-universal gaugino masses: Results

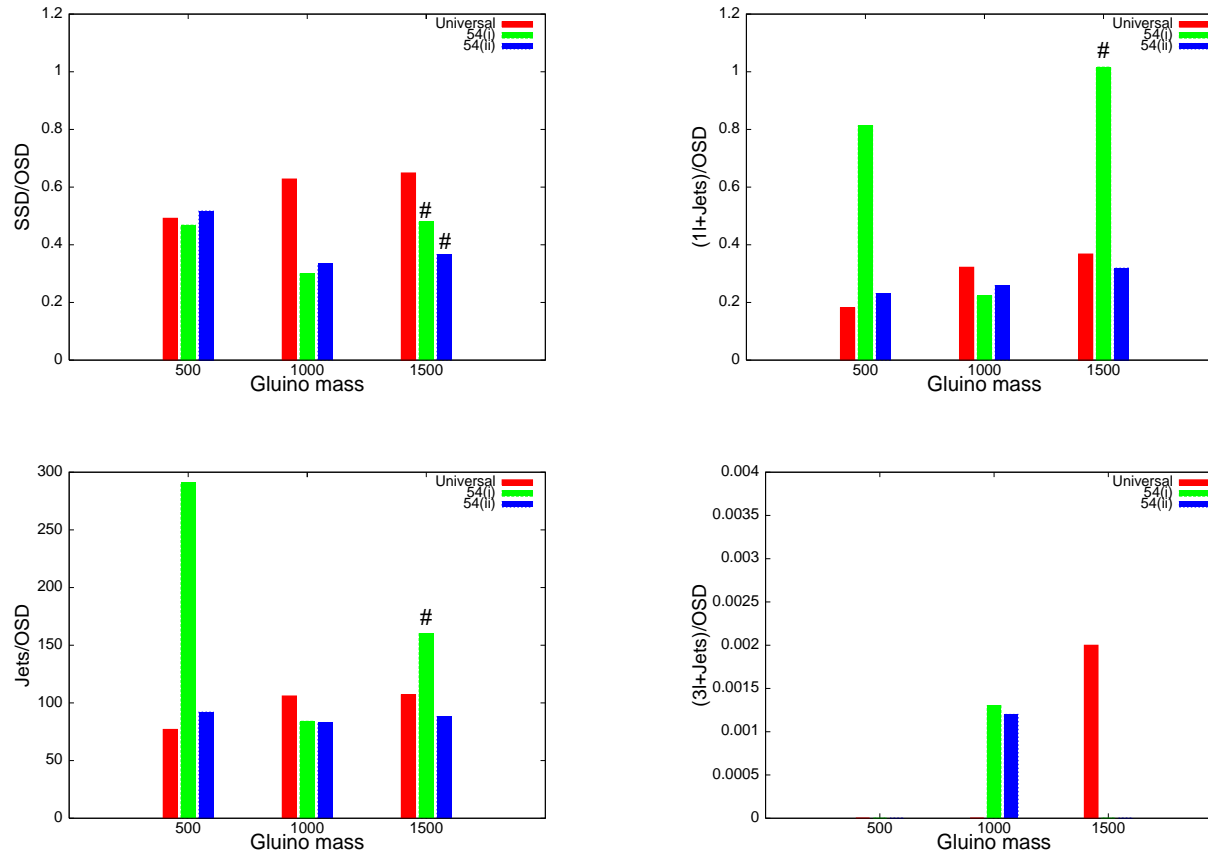


Figure 2: Event ratios for pMSSM in SO(10): $m_{\tilde{f}} = 1000$ GeV, $\mu = 1000$ GeV, $\tan \beta = 5$

Non-universal gaugino masses: Conclusions

- In a substantial region of the parameter space \longrightarrow **75 and 200** of SU(5) and **54 (i)** of SO(10) **easily distinguishable**.
- **24** of SU(5), **54(ii)** of SO(10) and the **universal case** \longrightarrow distinction is **relatively difficult**.
- **Trilepton channel is the most efficient discriminator**.
- Extraction of μ in pMSSM kind of framework is a challenging task. \longrightarrow **important**.

Non-universal scalar mass: Model1

- Model 1: Squark-Slepton Non-universality
 - Squarks and sleptons evolved from \longrightarrow mutually uncorrelated mass parameters $m_{0\tilde{q}}$ and $m_{0\tilde{l}}$ respectively.

Squark-slepton Non-universality: Results

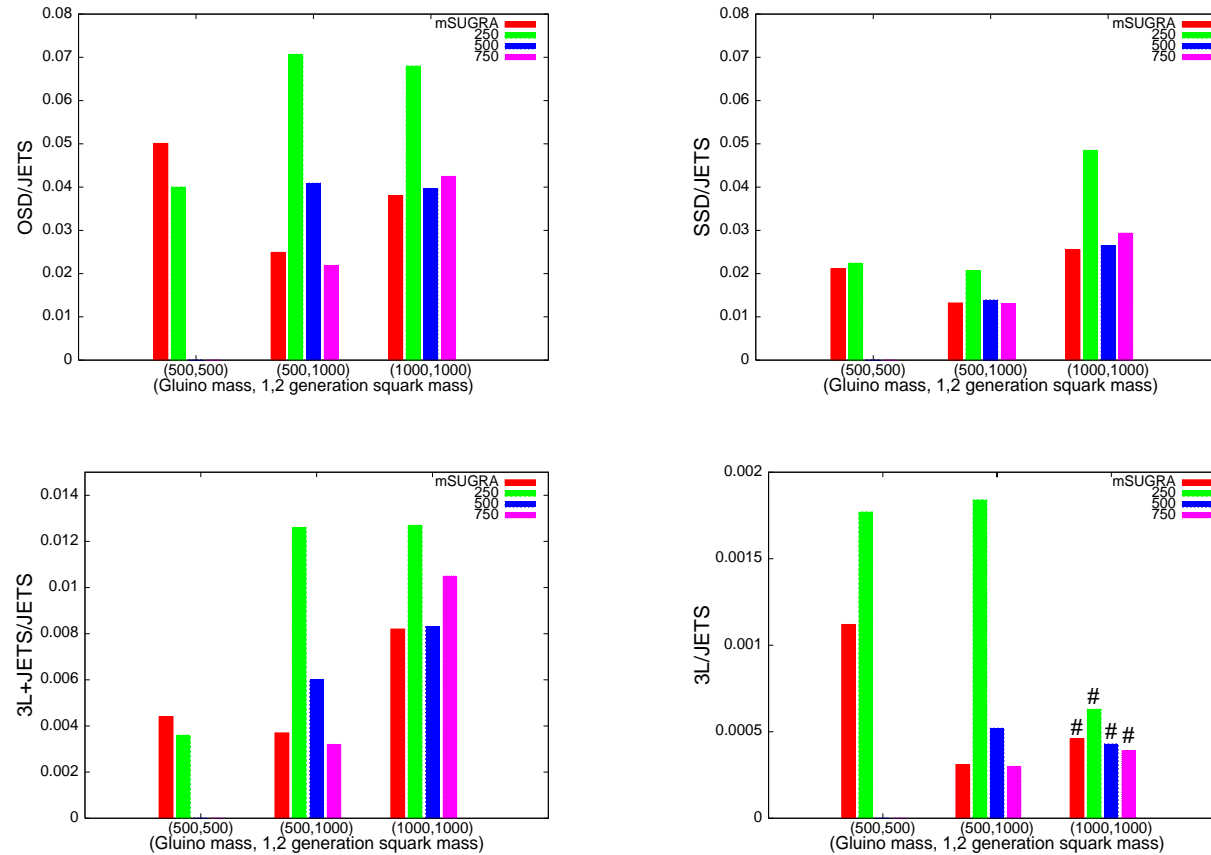


Figure 3: Event ratios for Squark-slepton Non-universality:
 $\tan \beta = 5$

Squark-slepton Non-universality: Results

- Cases with $m_{\tilde{1},2} = 250 \text{ GeV}$, is **fairly distinguishable** \longrightarrow especially for squark masses on the higher side.
- The $3\ell + jets$ events distinguish $m_{\tilde{1},2} = 750 \text{ GeV}$ \longrightarrow more prominent for high gluino mass and large $\tan\beta$.
- Cases with $m_{\tilde{1},2} = 500 \text{ GeV}$ \longrightarrow **difficult to differentiate from universal case.**

Non-universal scalar mass: Model2

- Model 2: Third family scalar non-universality
 - Third family scalars evolve from separate mass parameter m_0^3 from that of first two families $m_0^{(1,2)}$.
 - 1,2 families scalars may be very heavy \longrightarrow so called 'inverted hierarchy' \longrightarrow suppresses FCNC.

3rd family scalar Non-universality: Results

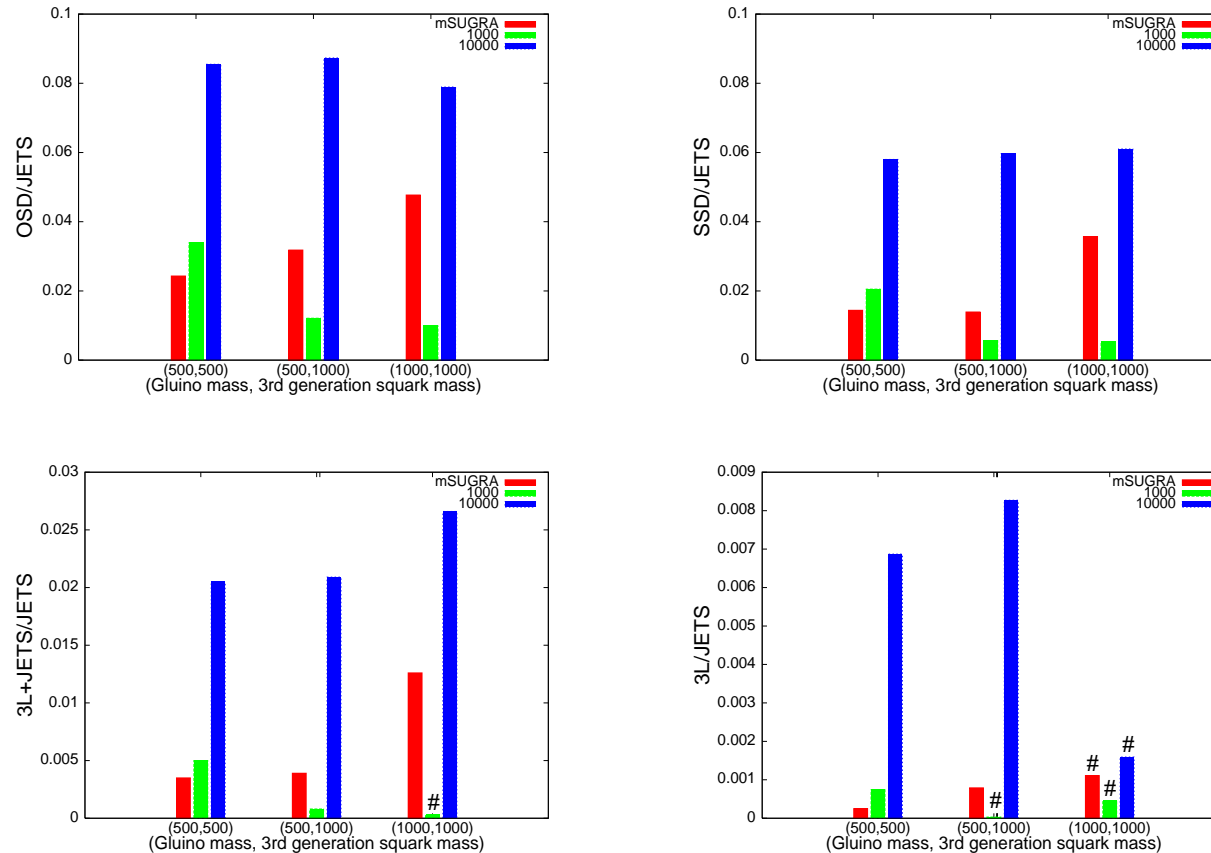


Figure 4: Event ratios for 3rd generation scalar Non-universality: $\tan \beta = 5$

3rd family scalar non-universality: Results

- The ratios are significantly higher for $m_{\tilde{q}^{1,2}} = 10$ TeV \longrightarrow Distinguishable.
- The ratios are significantly smaller for $m_{\tilde{q}^{1,2}} = 1$ TeV \longrightarrow Distinguishable.
- Unlike the other cases \longrightarrow very little dependence on the value of $\tan \beta$.

Non-universal scalar masses: Model3

- Model 3: Non-universality due to $SO(10)$ D -term
 - Matter fields belong to rep **16** \longrightarrow further classified into **submultiplets** \longrightarrow depending on the representations of $SU(5)$ to which they belong.
 - $\bar{5}(D^c \ \& \ L)$ or $10(E^c, U^c \ \& \ Q)$.
 - Breakdown of $SO(10)$ to SM gives \longrightarrow **different D -terms for different $SU(5)$ rep** .
 - Respectively for $\bar{5}$ and 10 :

$$m_{\bar{5}}^2 = m_0^2 - 1.5Dm_0^2 \quad (\text{for } D^c \ \& \ L)$$
$$m_{10}^2 = m_0^2 + 0.5Dm_0^2 \quad (\text{for } E^c, U^c \ \& \ Q)$$

Scalar Non-universality due to $SO(10)$ $D - term$: Results

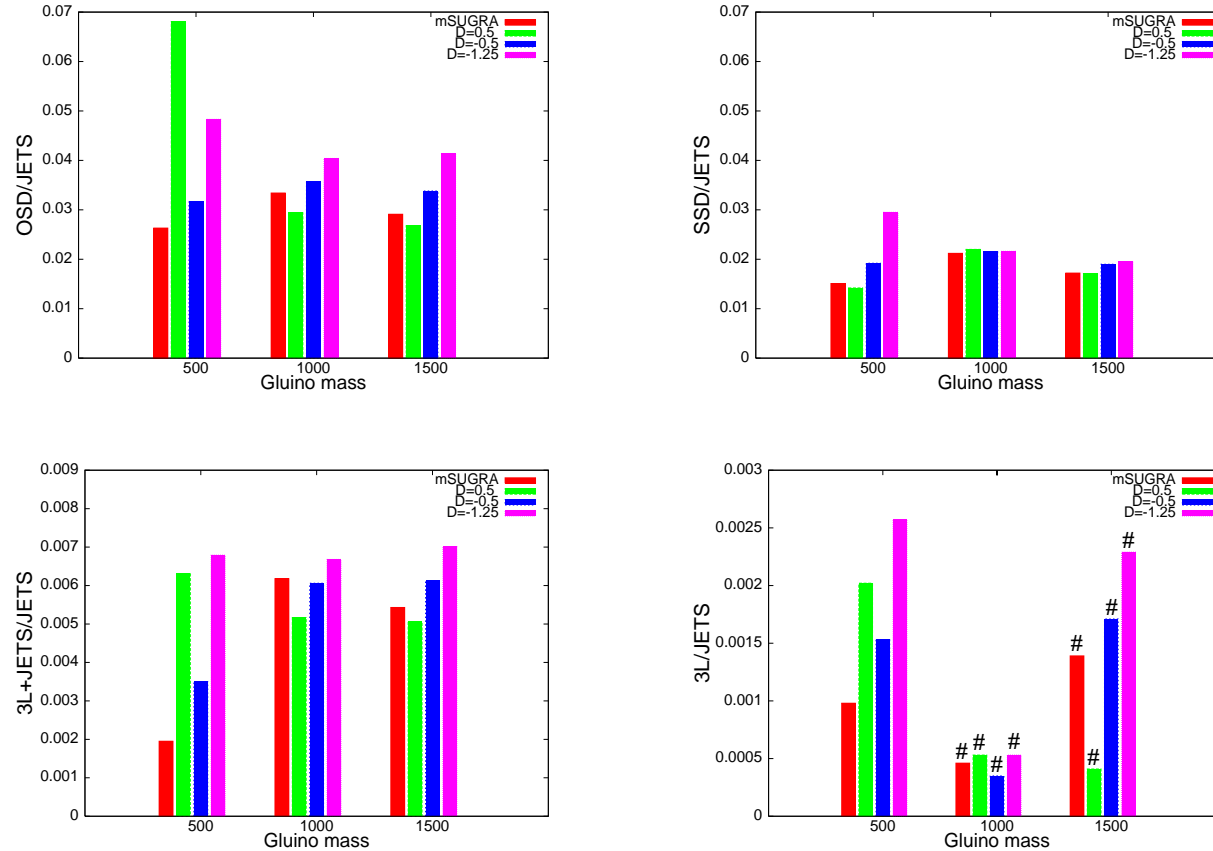


Figure 5: Event ratios for Scalar Non-universality due to $SO(10)$ $D - term$: $\tan \beta = 5$

Scalar Non-universality due to $SO(10)$ D – term : Results

- For $m_{\tilde{g}} = 1$ TeV and 1.5 TeV \longrightarrow distinction between $D = 0.5, -0.5$ and -1.25 \longrightarrow difficult from the ratio plot.
- For $m_{\tilde{g}} = 500$ GeV \longrightarrow $D = 0.5$ and $D = -1.25$ easily distinguishable from the ratios.
- The hadronically quiet trilepton \longrightarrow largely washed out by backgrounds excepting for $m_{\tilde{g}} = 500$ GeV.

Non-universal scalar mass: Conclusions

- Unlike gaugino non-universality \longrightarrow schemes of scalar non-universality **more non-uniform**.
- **Easiest identification** \longrightarrow **1,2 family very heavy** \longrightarrow 'inverted hierarchy'.
- Most difficult \longrightarrow **Various D-terms**, particularly for high gluino mass.

Thank You