

M-theory on pp-waves with a holomorphic superpotential and its matrix description

Nakwoo Kim

Physics Department
Kyung Hee University

SUSY 08, COEX Seoul Korea, 20 Apr 2008

Based on a work with J. Kim, J.-H. Park, J. Plefka arXiv:0804.3349

Matrix theory conjecture

- The discovery of D-branes in string theory among other things suggests that Yang-Mills theory can give a nonperturbative formulation of Quantum Gravity.
- Concrete examples are Matrix theory : BFSS for M-theory, Matrix strings for IIA, IKKT for IIB etc. item AdS/CFT correspondence is the most elaborate and well-studied : 4d SCFT and Strings in Curved spacetime.

Examples and Constructions of Matrix Theory

- Take M-theory in 11d, compactify on a circle to get IIA, and give infinite boost along x^{11} .
 - Infinite momentum becomes infinite number of D0-branes in IIA, and the conjecture is that the D0-brane dynamics in the large N limit gives an M-theory description.
 - Alternatively one can consider supermembrane action choose light-cone gauge and discretize the Poisson bracket as matrix commutator.
- The result is maximally supersymmetric $U(N)$ Yang-Mills quantum mechanics.
- Successfully describes M2, M5, scattering amplitudes etc.

Pp-waves

- Unfortunately the Matrix theory conjecture is very much background-dependent but can be applied to any pp-waves.
- Generalized to max. susy pp-wave by [Berenstein, Maldacena, Nastase \(JHEP 0204, 013 2002\)](#) : Shows several attractive features like possibility of perturbative computation of energy spectrum, transverse M5-branes etc.

Inhomogeneous pp-waves

- We are here interested in pp-waves with **non-constant** flux.

$$ds^2 = 2dx^+ dx^- + H(x^+, x^M)(dx^+)^2 + \sum_{M=1}^9 (dx^M)^2$$

$$G = dx^+ \wedge \phi(x^+, x^M)$$

- Einstein equation becomes $\nabla^2 H = -\phi^2/6$. ϕ is closed and co-closed.
- We implement **6+3 split of transverse space, and allow coordinate dependence on 6d only**. (which is to be complexified.)

Pp-wave with holomorphic superpotential

- All the equations are satisfied if we set

$$\phi_{\bar{a}bc} = \partial_{\bar{a}} \partial_{\bar{d}} \bar{W} \epsilon^{\bar{d}}_{bc}, \quad H = -|\partial W|^2$$

- It can be verified that there exist in general 4 kinematical and 4 dynamical susy.
- SUSY projection : $\Gamma_{12}\epsilon = \Gamma_{34}\epsilon = \Gamma_{56}\epsilon$: 1/4-BPS

Supermembrane action

Use the 1st order formulation and choose the light-cone gauge, then we have - in terms of **Poisson bracket**.

$$\mathcal{L} = \frac{1}{2}(D_t X_M)^2 + \frac{1}{2}H + \frac{1}{2}C_{+rs}\{X^r, X^s\} - \frac{1}{4}\{X^M, X^N\}^2 + \text{fermions}$$

Now use the form of our solution, and express in terms of 3 complex coordinates instead of 6 real (3 remains the same)

$$\begin{aligned} \mathcal{L} = & \frac{1}{2}(D_t X_M)^2 - \frac{1}{4}\{X^M, X^N\}\{X_M, X_N\} + i\theta^\dagger D_t \theta + i\theta^\dagger \Gamma_M \{X^M, \theta\} \\ & - \frac{1}{2}\partial_a W \partial_{\bar{a}} \bar{W} + \frac{1}{2}\epsilon_{dab}\partial_{\bar{d}} \bar{W} \{Z^a, Z^b\} + \frac{1}{2}\epsilon_{\bar{d}\bar{a}\bar{b}}\partial_d W \{Z^{\bar{a}}, Z^{\bar{b}}\} \\ & - \frac{i}{8}\epsilon_{\bar{d}\bar{a}\bar{b}}\partial_c \partial_d W \theta^\dagger \Gamma_{\bar{c}ab} \theta - \frac{i}{8}\epsilon_{dab}\partial_{\bar{c}} \partial_{\bar{d}} W \theta^\dagger \Gamma_{c\bar{a}\bar{b}} \theta \end{aligned}$$

Discretization and Gauge Field theory

- Now we can follow the standard discretization procedure and replay the coordinates by $N \times N$ **matrices**. Poisson bracket becomes matrix commutators.
- The 6+3 split, $SO(3)$ symmetry, and the solution being described in terms of a holomorphic function, all suggest that things are naturally related to $N = 1$ supersymmetric theory in 3+1 dimensions.
- It is indeed quite convenient to view the action as a dimensional reduction of $D = 4$, $N = 1$ theory to quantum mechanics.

Matrix Model Action

$$\begin{aligned} \mathcal{L} = & \text{Tr} \left(\frac{1}{2} D_t X^M D_t X_M + \frac{1}{4} [X^M, X^N] [X_M, X_N] + i \theta^\dagger D_t \theta - \theta^\dagger \Gamma^M [X_M, \theta] \right) \\ & + \frac{1}{2} \text{Tr} \left(i \epsilon_{\bar{a}\bar{b}}{}^c [\bar{Z}^{\bar{a}}, \bar{Z}^{\bar{b}}] \partial_c W + i \epsilon_{ab}{}^{\bar{c}} [Z^a, Z^b] \bar{\partial}_{\bar{c}} \bar{W} - \partial_a W \partial^a \bar{W} \right) \\ & - \frac{i}{8} \text{Tr} \left(\theta^\dagger \Gamma^a \Gamma^{\bar{1}\bar{2}\bar{3}} \Gamma^b \partial_a \text{Tr} (\theta \partial_b W) + \theta^\dagger \Gamma^{\bar{a}} \Gamma^{123} \Gamma^{\bar{b}} \bar{\partial}_{\bar{a}} \text{Tr} (\theta \bar{\partial}_{\bar{b}} \bar{W}) \right) \end{aligned}$$

- Goes over to the membrane action in $N \rightarrow \infty$ limit.
- W is now (multi)-trace, holomorphic operator written using 3 complex matrices.
- Exists ambiguity to get a (matrix) W from supergravity W .
- We can use $\mathcal{W} = W - \frac{i}{3} \epsilon_{abc} \text{Tr} (Z^a, [Z^b, Z^c])$ for the whole superpotential.

Discussions

- Obtained an interesting class of inhomogeneous pp-waves in 11d and found the matrix action.
- The relation between 11d pp-wave and gauge theory turns out to be very rich.
- Might concentrate on more interesting backgrounds :
 β -deformation of BFSS (addressed by H. Shimada)
- One can try to generalize to other dimensions. (7+2 or 8+1 split).