

1-Loop Effective Action on Orbifolds

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Introduction

- Orbifolds are a very popular choice for the *internal geometry* for String Theory and Field Theory with Extra Dimensions
- Simplest Geometry allowing 4d chiral fermions, and $N = 1$ SUSY
- Plethora of *Symmetry Breaking* Possibilities
- Particular limits of more general spaces (CY)
- High degree of *calculability!*

Introduction

- Special Interest in *perturbative corrections*
- *Moduli stabilization* problem: Many 4d scalars with no potential at tree level
 - No-Scale SUSY breaking
- Often: Interested in *Large Volume* →
Perturbation theory both reliable and dominant over nonperturbative physics....
- Want to go beyond eff. potential approximation

The covariant background method

[DeWitt 1965; Gilkey 1975]

- Split all fields in background and fluctuations

$$\phi(x) = \varphi(x) + \varphi_{dyn}(x)$$

$$A_M(x) = A_M(x) + A_{dyn,M}(x) \quad \text{etc} \dots$$

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Heat Kernel $K(x, x', T)$



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$$D_M = \nabla_M - iA_M - i\omega_M$$

$$E = E(\varphi, A_M, g_{MN})$$

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- Describes effective action including operators with derivatives
- Effective potential is obtained by setting to zero all derivatives...

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- Nontrivial λ gives nonlocal renorm. (finite)

...Torus...

- Basic idea: Further expand heat kernel in powers of λ : [GG 08]

$$K(x,x,T) = \frac{\sqrt{g}}{(4\pi T)^{d/2}} \left(\sum_r a_r T^r \right)$$

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$$K(x, x-\lambda, T) = \frac{\sqrt{g}}{(4\pi T)^{d/2}} \exp\left(\frac{-\lambda^2}{4T}\right) \left(\sum_{r,s} a_{r, j_1 \dots j_s} T^r \lambda^{j_1} \dots \lambda^{j_s} \right) \exp(i\lambda^j [A_j + \omega_j])$$

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[GG 08]

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- Exponential prefactor: Suppression for nonzero λ , finite (nonlocal) contributions
- Wilson line \rightarrow partial resummation
- Has been calculated up to dim 4 operators

Result up to dim 4 operators

$$K(x, x - \lambda, T) = \frac{\sqrt{g}}{(4\pi T)^{d/2}} \exp\left(-\frac{\lambda^2}{4T}\right) \quad [\text{GG 08}]$$

$$\left\{ 1 + T \left[\frac{1}{6} R - E \right] + T \lambda^j \left[-\frac{1}{12} R_{;j} + \frac{1}{2} E_{;j} - \frac{1}{6} F^M_{j;M} \right] \right.$$

$$+ T \lambda^j \lambda^k \left[\frac{1}{40} R_{;jk} + \frac{1}{120} R_{jk;M}^M - \frac{1}{90} R^M_j R_{Mk} \right.$$

$$+ \frac{1}{180} R^{MN} R_{MjNk} + \frac{1}{180} R^{MNL}_j R_{MNLk} - \frac{1}{6} E_{;jk}$$

$$\left. + \frac{1}{24} R^M_j F_{Mk} + \frac{1}{12} F_{Mj} F^M_k + \frac{1}{12} F_{Mj;M}^k \right]$$

$$+ T^2 \left[\frac{1}{2} \left(\frac{1}{6} R - E \right)^2 + \frac{1}{6} \left(\frac{1}{5} R - E \right)_{;M}^M - \frac{1}{180} R_{MN} R^{MN} \right.$$

$$\left. + \frac{1}{180} R^{MNLS} R_{MNLS} + \frac{1}{12} F_{MN} F^{MN} \right] \left. \right\} \exp(i\lambda^j [A_j + \omega_j])$$

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- Contributions with $\lambda = 0$ now correspond to renormalization at fixed points of P^k
- Remains to evaluate $K(x, Px - \lambda)$

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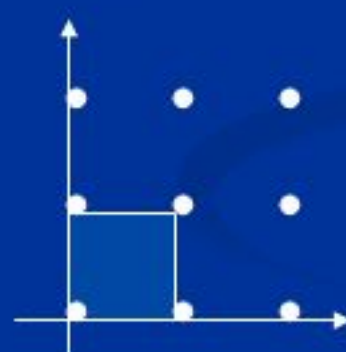
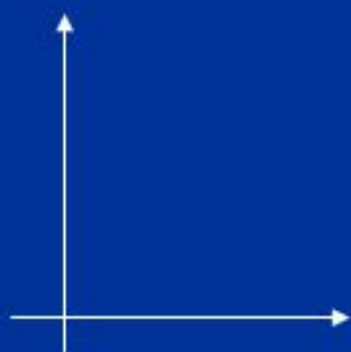
- For each orbifold sector k , split into fixed and rotated dimensions

\mathbb{R}^4 fixed

Λ_{\parallel} fixed

Λ_{\perp} rotated

Under P^k :



$$x = x_{\parallel} + x_{\perp}, \quad \Lambda = \Lambda_{\parallel} + \Lambda_{\perp}$$

- Either torus can be trivial for certain k

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$$K_{\parallel}(T) = \exp(-T[-D_{\parallel}^2 + E + A_{\perp}^2 + \omega_{\perp}^2])$$

$$\varphi = \varphi(x^{\mu})$$

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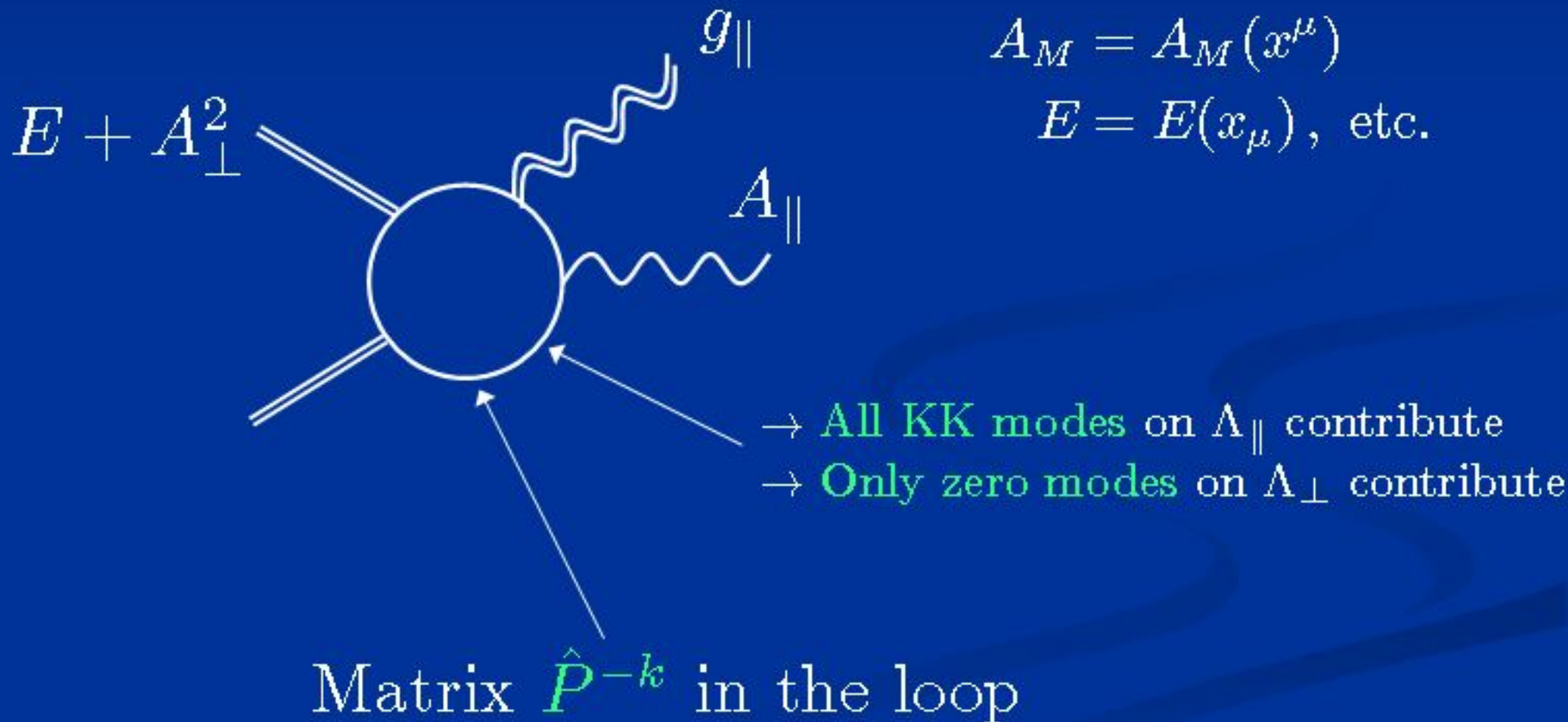
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Evaluation as on the d_{\parallel} -dimensional fixed torus with shifted mass matrix

Diagrammatic Interpretation



Discrete Wilson lines

- Fields need not be periodic on the torus

$$\phi(x + \lambda) = W_0(\lambda)\phi(x)$$

- Certain discrete WL cannot be transformed into backgrounds for gauge fields: $W_0(\lambda)^N = 1$

- Simple modifications

- Cts. WL is rescaled: $e^{i\lambda_{\parallel} \cdot A} \rightarrow e^{i\lambda_{\parallel} \cdot A} W_0(\lambda_{\parallel})$

- Additional matrix in trace: $\text{tr } K_{\parallel} \hat{P}^{-k} \rightarrow \text{tr } K_{\parallel} \hat{P}^{-k} Q$

Q is the projector onto zero modes on the transverse torus Λ_{\perp} carrying the discrete Wilson line $W_0(\lambda_{\perp})$

Conclusions

- Applied the covariant background (heat kernel) formalism to calculate 1-loop corrections on arbitrary torus orbifolds
- Result on torus: Covariant expansion in powers of lattice vectors, calculated all dim 4 operators
- Orbifold: Each sector corresponds to the result of its fixed torus with a shifted mass matrix