

# EWSB from Unparticles

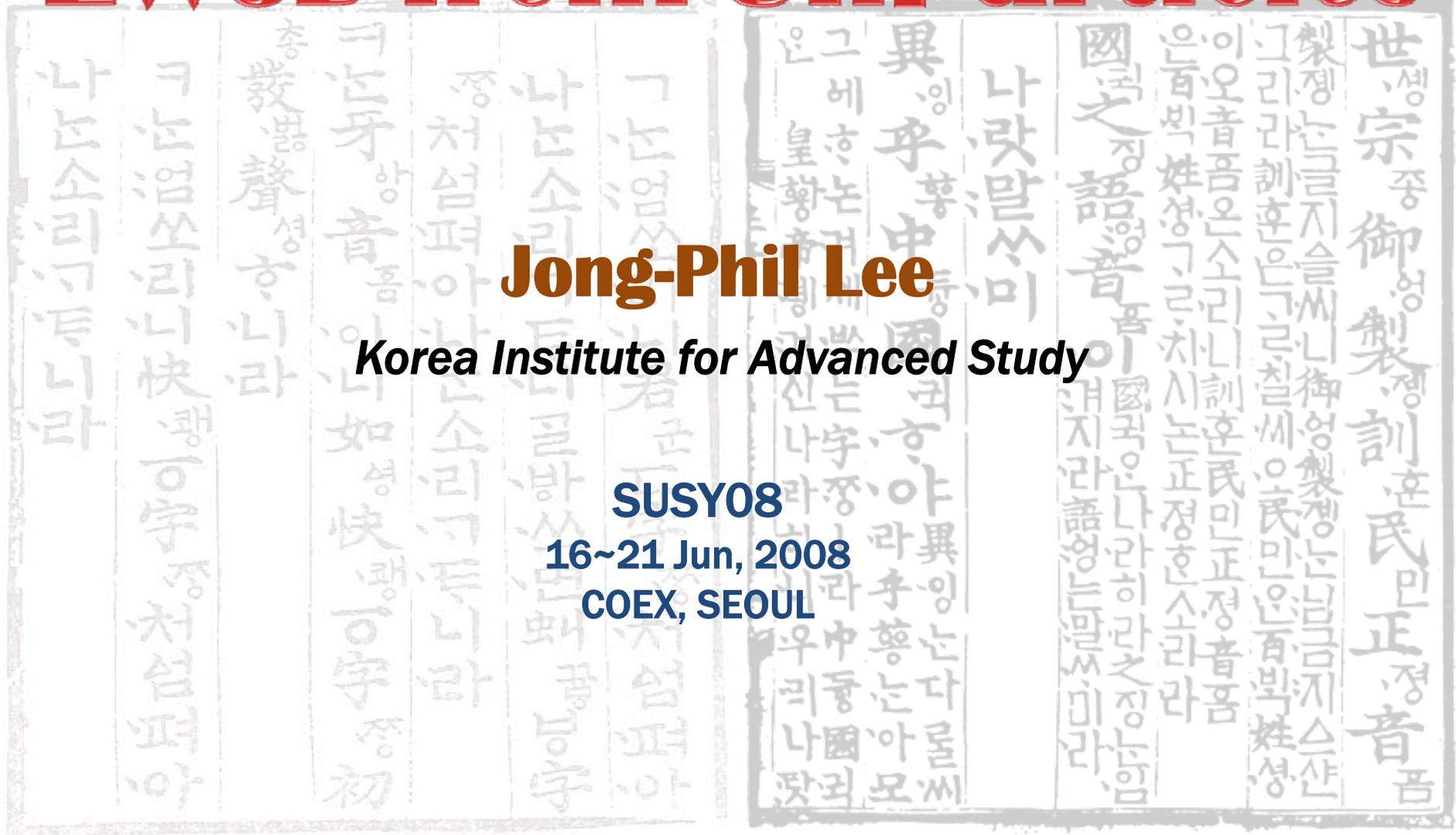
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**Korea Institute for Advanced Study**

**SUSY08**

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**COEX, SEOUL**



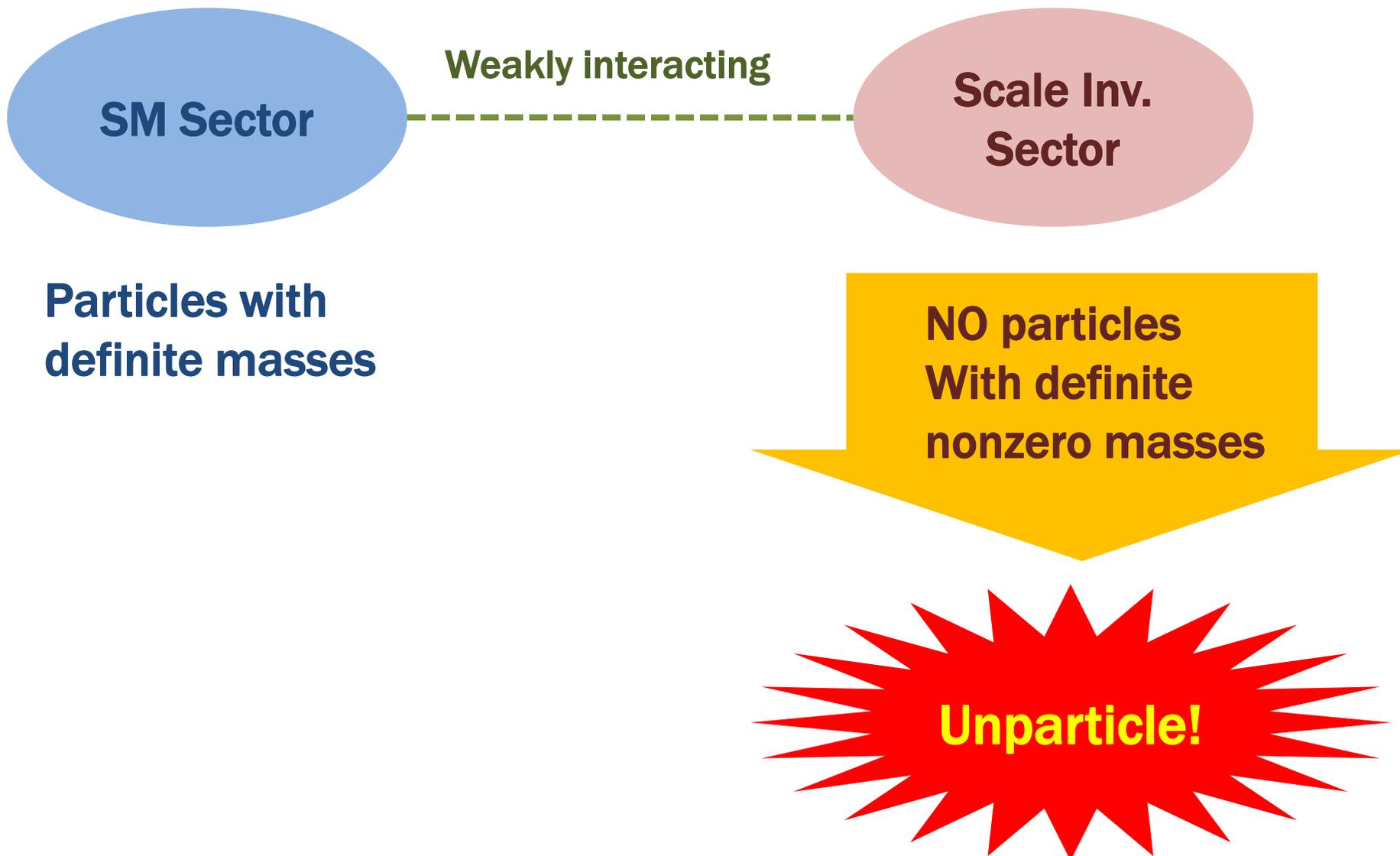
# Outlook

- Unparticle Brief
- New scalar potential
- EWSB with scalar potential
- Mass spectrum
- Conclusions

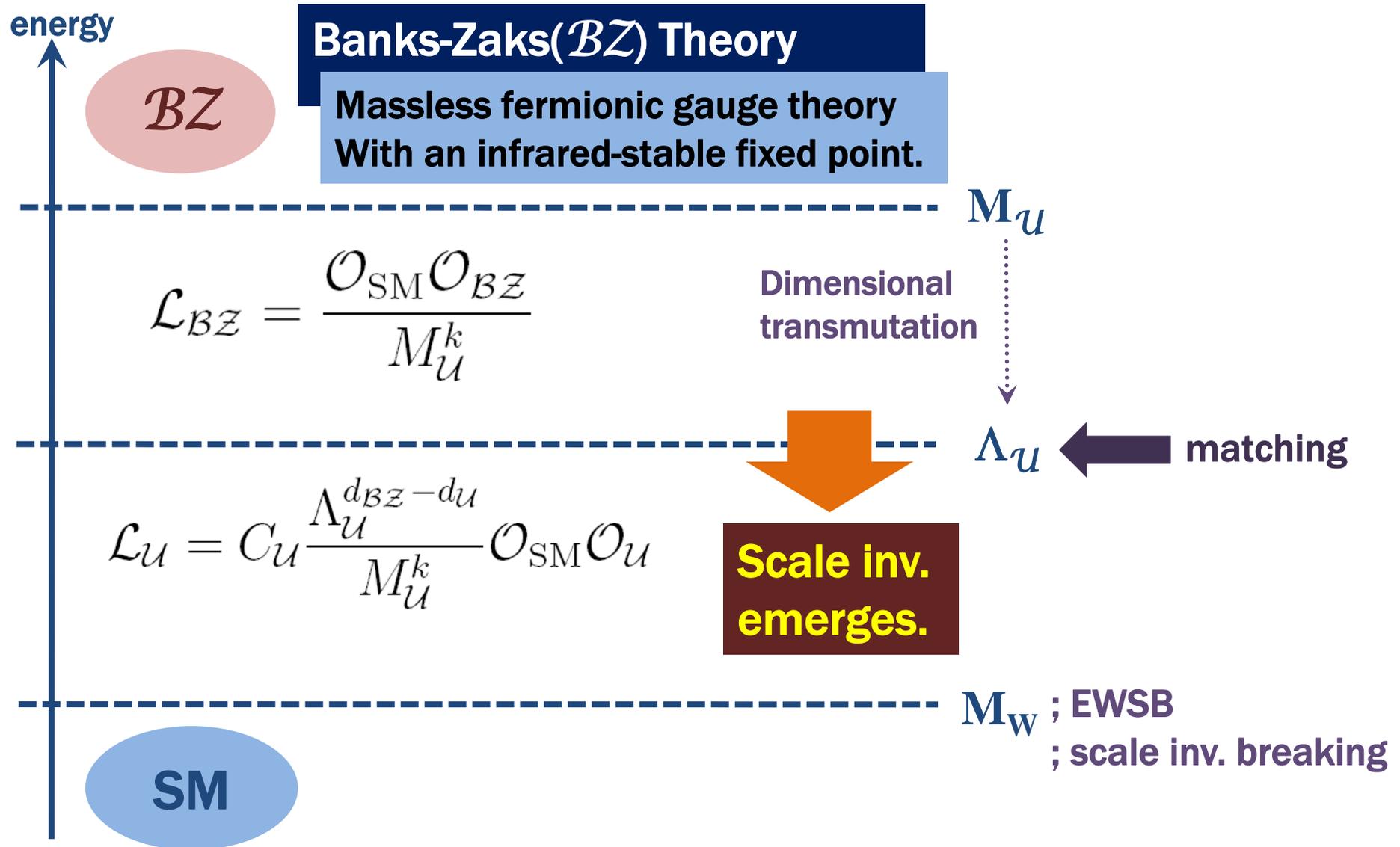
# Unparticles

# Unparticles $\mathcal{U}$ : Basic Idea

H. Georgi, PRL98; PLB650



# Effective Theory for $\mathcal{U}$



# Phase Space of $\mathcal{U}$

## Two-point function

$$\langle 0 | \mathcal{O}_{\mathcal{U}}(x) \mathcal{O}^{\dagger}(0) | 0 \rangle = \int \frac{d^4 P}{(2\pi)^4} e^{-iP \cdot x} \rho_{\mathcal{U}}(P^2)$$

## Spectral density function

$$\rho_{\mathcal{U}}(P^2) = A_{d_{\mathcal{U}}} \theta(P^0) \theta(P^2) (P^2)^{d_{\mathcal{U}}-2}$$

Fixed by scale inv.

## Normalization factor

$$A_{d_{\mathcal{U}}} = \frac{16\pi^{5/2}}{(2\pi)^{2d_{\mathcal{U}}}} \frac{\Gamma(d_{\mathcal{U}} + \frac{1}{2})}{\Gamma(d_{\mathcal{U}} - 1)\Gamma(2d_{\mathcal{U}})}$$

Unparticles with  $d_{\mathcal{U}}$  looks like a Non-integral number of massless particles.

## Scalar Unparticle Propagator

$$\Delta_F(P^2) = \frac{A_{d\mathcal{U}}}{2 \sin(d\mathcal{U}\pi)} (-P^2)^{d\mathcal{U}-2}$$

## Vector Unparticle Propagator

$$\int d^4x e^{-ik \cdot x} \langle 0 | T(O_\mu(x) O_\nu(0)) | 0 \rangle$$
$$= -iC (-k^2 - i\epsilon)^{d-3} \left[ k^2 g_{\mu\nu} - \frac{2(d-2)}{d-1} k_\mu k_\nu \right]$$

# New Scalar Potential

## Relevant Interaction

$$\mathcal{L}_{\Phi\Phi\mathcal{U}} \sim \lambda_{\Phi\Phi\mathcal{U}}(\Phi^\dagger\Phi)\mathcal{O}_{\mathcal{U}}, \quad [\lambda_{\Phi\Phi\mathcal{U}}] = 2 - d_{\mathcal{U}} > 0$$

## New proposal

$$V(\Phi, \phi) = \lambda_0(\Phi^\dagger\Phi)^2 + \lambda_1(\phi^*\phi)^2 + 2\lambda_2\mu^{2-d_{\mathcal{U}}}(\Phi^\dagger\Phi)(\phi^*\phi)^{d_{\mathcal{U}}/2}$$

SM scalar
Hidden scalar
New interaction

# ElectroWeak Symmetry Breaking

# EWSB with Scalar Potential

## Ordinary Higgs mechanism

$$V = \frac{m^2}{2} h^2 + \frac{\lambda}{4} h^4$$

## Coleman-Weinberg mechanism

$$V = \frac{\lambda}{4!} \varphi_c^4 + \frac{\lambda^2 \varphi_c^4}{256\pi^2} \left( \ln \frac{\varphi_c^2}{M^2} - \frac{25}{6} \right)$$

## Gildener-Weinberg mechanism

$$V_0(\Phi) = \frac{1}{24} f_{ijkl} \Phi_i \Phi_j \Phi_k \Phi_l.$$

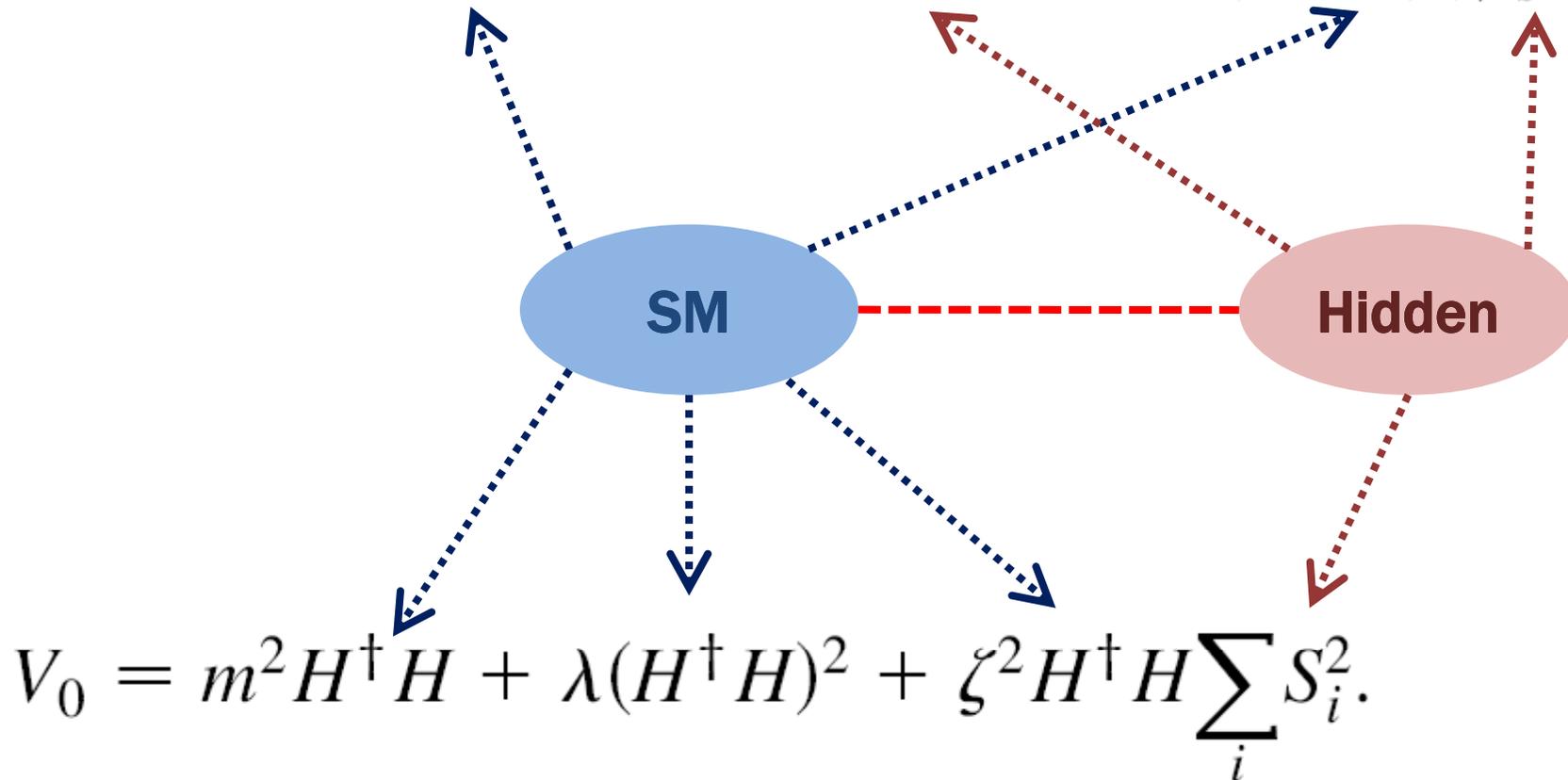
$$\delta V(n\phi) = A\phi^4 + B\phi^4 \ln(\phi^2/\Lambda_W^2)$$

$$A = \frac{1}{64\pi^2 \langle \phi \rangle^4} \{ 3 \text{Tr}[\mu^4 \ln(\mu^2/\langle \phi \rangle^2)] \\ + \text{Tr}[M^4 \ln(M^2/\langle \phi \rangle^2)] \\ - 4 \text{Tr}[m^4 \ln(m^2/\langle \phi \rangle^2)] \}$$
$$B = \frac{1}{64\pi^2 \langle \phi \rangle^4} (3 \text{Tr} \mu^4 + \text{Tr} M^4 - 4 \text{Tr} m^4)$$

# EWSB from Hidden Sector

Chang, Ng, Wu, PRD75;  
Espinosa, Quiros, PRD76

$$V_0(\Phi, \phi_s) = \lambda(\Phi^\dagger \Phi)^2 + \lambda_s(\phi_s^* \phi_s)^2 + 2\kappa(\Phi^\dagger \Phi)(\phi_s^* \phi_s)$$



# Unparticle-Higgs Interplay

## SM potential

$$V_0 = m^2 |H|^2 + \lambda |H|^4$$

## Deconstruction

$$\mathcal{O} \equiv \sum_n F_n \varphi_n$$

$$F_n^2 = \frac{A_{d_U}}{2\pi} \Delta^2 (M_n^2)^{d_U - 2}$$

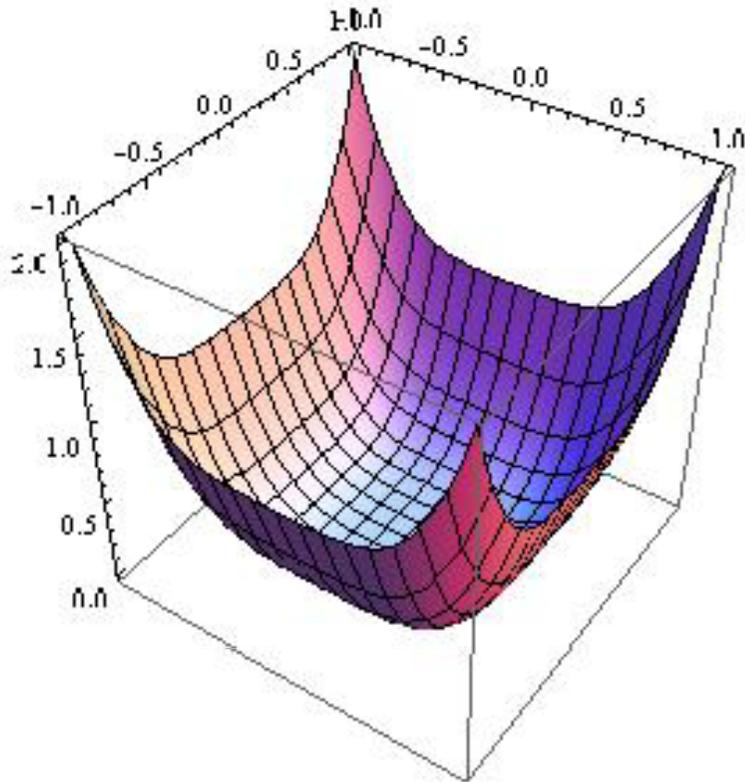
## U-sector contribution

$$\delta V = \frac{1}{2} \sum_n M_n^2 \varphi_n^2 + \kappa_U |H|^2 \sum_n F_n \varphi_n$$

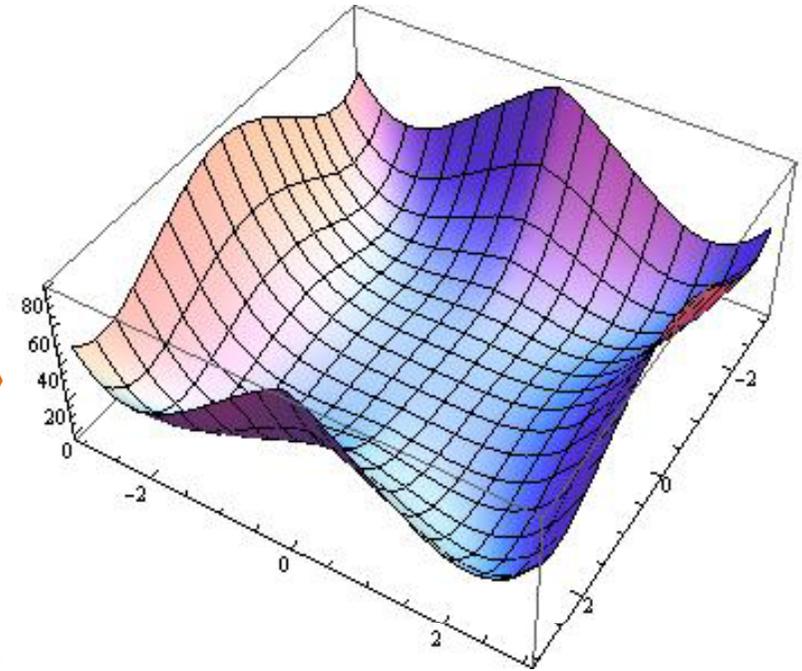
## IR regulator

$$\delta V = \zeta |H|^2 \sum_n \varphi_n^2$$

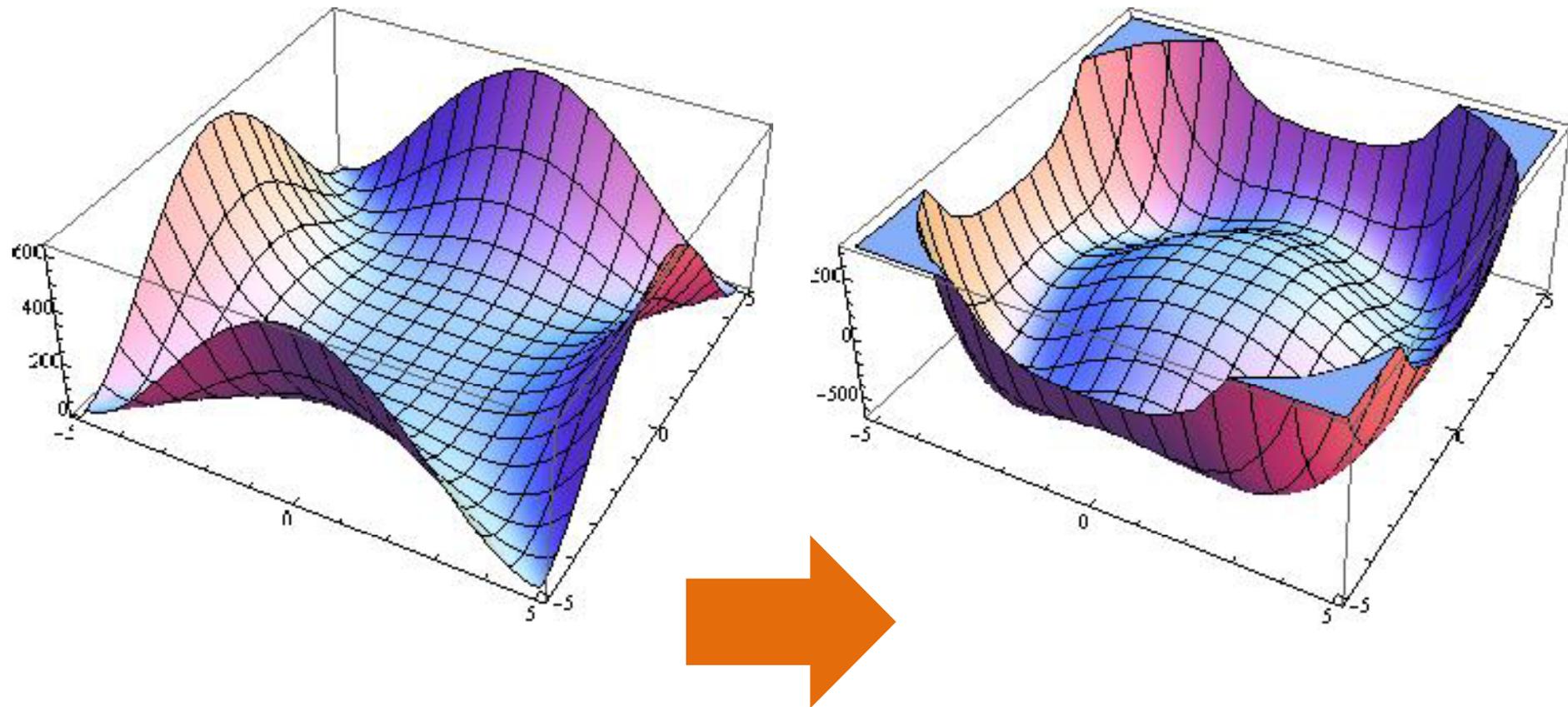
# Examples:



**EWSB**



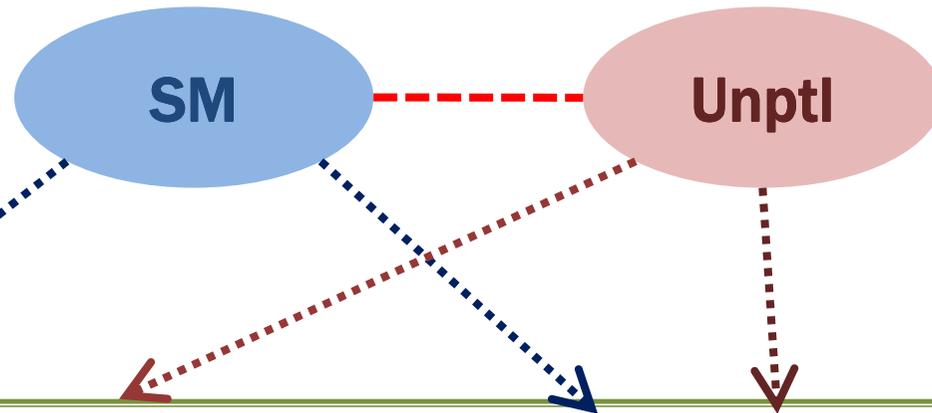
# Gildener-Weinberg Mechanism



**Coleman-Weinberg turned on**

# Finding Minimum of New Potential

## New Potential



$$V(\Phi, \phi) = \lambda_0(\Phi^\dagger \Phi)^2 + \lambda_1(\phi^* \phi)^2 + 2\lambda_2 \mu^{2-d_U} (\Phi^\dagger \Phi)(\phi^* \phi)^{d_U/2}$$

## Stationary Condition

$$(\partial V / \partial N_i)_{\vec{n}} = 0$$

$$\left. \frac{\partial^2 V}{\partial N_i \partial N_j} \right|_{\vec{n}} u_i u_j \geq 0$$

$$\Phi_i = (\Phi, \phi) = \rho N_i$$

Find the minimum  
along **this ray**  
for unit vector  $N_i$

# EWSB with New Potential

Stationary Condition satisfied when

$$\left(\frac{\hat{\rho}^2}{2}\right)^{-\epsilon} \lambda_2 n_1^{d_U} = -\lambda_0 n_0^2 \quad (d_U = 2 - 2\epsilon)$$
$$2\lambda_1 n_1^4 = d_U \lambda_0 n_0^4$$

Unit Vectors

$$n_0^2 = \frac{\sqrt{2\lambda_1}}{\sqrt{d_U \lambda_0} + \sqrt{2\lambda_1}} \quad n_1^2 = \frac{\sqrt{d_U \lambda_0}}{\sqrt{d_U \lambda_0} + \sqrt{2\lambda_1}}$$

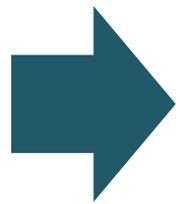
# VEV of $\rho$

$\rho_0$  is fixed:

$$\rho = \rho_0 \equiv \left( -\frac{2^\epsilon \lambda_2 n_1^{d_u}}{\lambda_0 n_0^2} \right)^{\frac{1}{2\epsilon}} \mu$$

Stability condition for  $\rho_0$

$$\lambda_2 = -\sqrt{\lambda_0 \lambda_1} \equiv \bar{\lambda}$$



$$\hat{\rho}_0^2 = 2 \left( \frac{d_u}{2} \right)^{\frac{1}{2-d_u}} \frac{\sqrt{d_u \lambda_0} + \sqrt{2\lambda_1}}{\sqrt{d_u \lambda_0}}$$
$$\rightarrow \frac{2}{\sqrt{e}} \frac{\sqrt{\lambda_0} + \sqrt{\lambda_1}}{\sqrt{\lambda_0}}, \quad \text{as } d_u \rightarrow 2$$

# Constraints

From the weak boson mass:

$$(n_0 \rho_0)^2 = \frac{1}{\sqrt{2} G_F} = (246 \text{ GeV})^2 \equiv v_0^2$$

If we choose  $\mu = v_0$ :

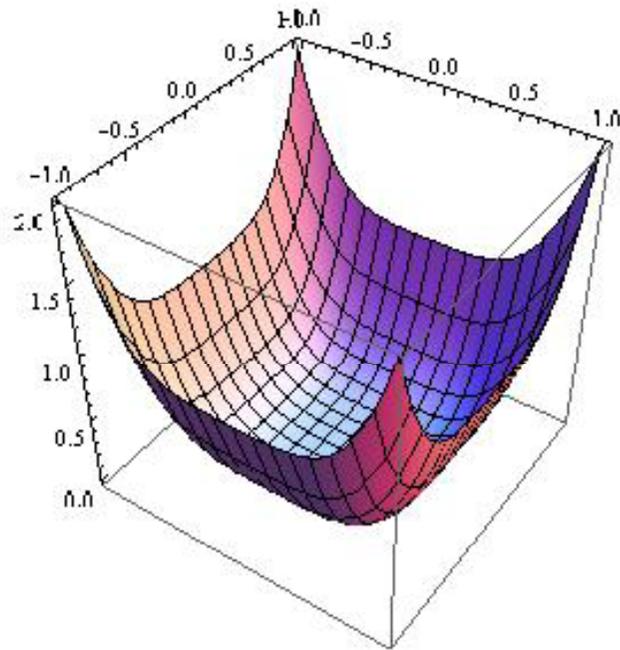
$$\frac{\lambda_1}{\lambda_0} = \frac{1}{4} \left( \frac{2}{d_{\mathcal{U}}} \right)^{\frac{d_{\mathcal{U}}}{2-d_{\mathcal{U}}}}$$

$$\longrightarrow \frac{e}{4} \simeq 0.68 \quad \text{as } d_{\mathcal{U}} \rightarrow 2$$

# How to break EWS

SM

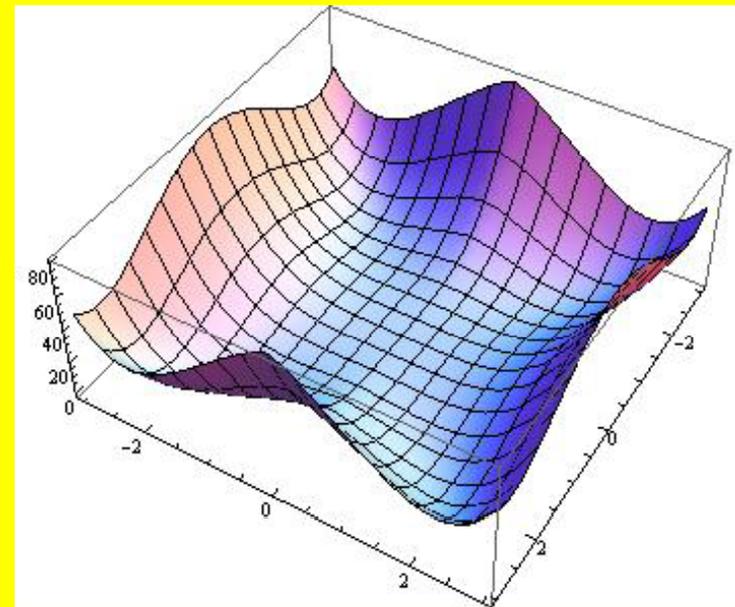
Scale  
Inv.



$$f(x,y) = x^4 + y^4$$

SM

Scale  
Inv.



$$f(x,y) = x^4 + y^4 - 3x^2y^{1.3}$$

# Mass Spectrum

# Fluctuations around the vacuum

## Fields redefined

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ n_0 \rho_0 + h \end{pmatrix}, \quad \phi = \frac{1}{\sqrt{2}} (n_1 \rho_0 + s)$$

## Potential of $h$ and $s$

$$V(h, s) = \frac{\lambda_0}{4} (n_0 \rho_0 + h)^4 + \frac{\lambda_1}{4} (n_1 \rho_0 + s)^4 \\ + 2^{-d\mu/2} \lambda_2 \mu^{2\epsilon} (n_0 \rho_0 + h)^2 (n_1 \rho_0 + s)^{d\mu}$$

# Mass Squared Matrix

$\psi_i = (h, s)$

$$(M^2)_{i,j} = \left. \frac{\partial^2 V}{\partial \psi_i \partial \psi_j} \right|_0$$

$$= \frac{\rho_0^2 n_0^2}{\sqrt{2\lambda_1}} \begin{pmatrix} 2\lambda_0 \sqrt{2\lambda_1} & - (2d_{\mathcal{U}} \lambda_0 \lambda_1)^{\frac{3}{4}} \\ - (2d_{\mathcal{U}} \lambda_0 \lambda_1)^{\frac{3}{4}} & (4 - d_{\mathcal{U}}) \lambda_1 \sqrt{d_{\mathcal{U}} \lambda_0} \end{pmatrix}$$

# Mass Spectrum

$$m_{h,\ell}^2 = \frac{\rho_0^2 \sqrt{2\lambda_0\lambda_1}}{\sqrt{d_u\lambda_0} + \sqrt{2\lambda_1}} \left\{ \sqrt{\lambda_0} + \left(2 - \frac{d_u}{2}\right) \sqrt{\frac{d_u}{2}\lambda_1} \pm \sqrt{D} \right\}$$

$$D = \lambda_0 + \left(2 - \frac{d_u}{2}\right)^2 \frac{d_u}{2} \lambda_1 + \left(\frac{3d_u}{2} - 2\right) \sqrt{2d_u\lambda_0\lambda_1}$$

$$\epsilon = 1 - d_u/2 \ll 1$$

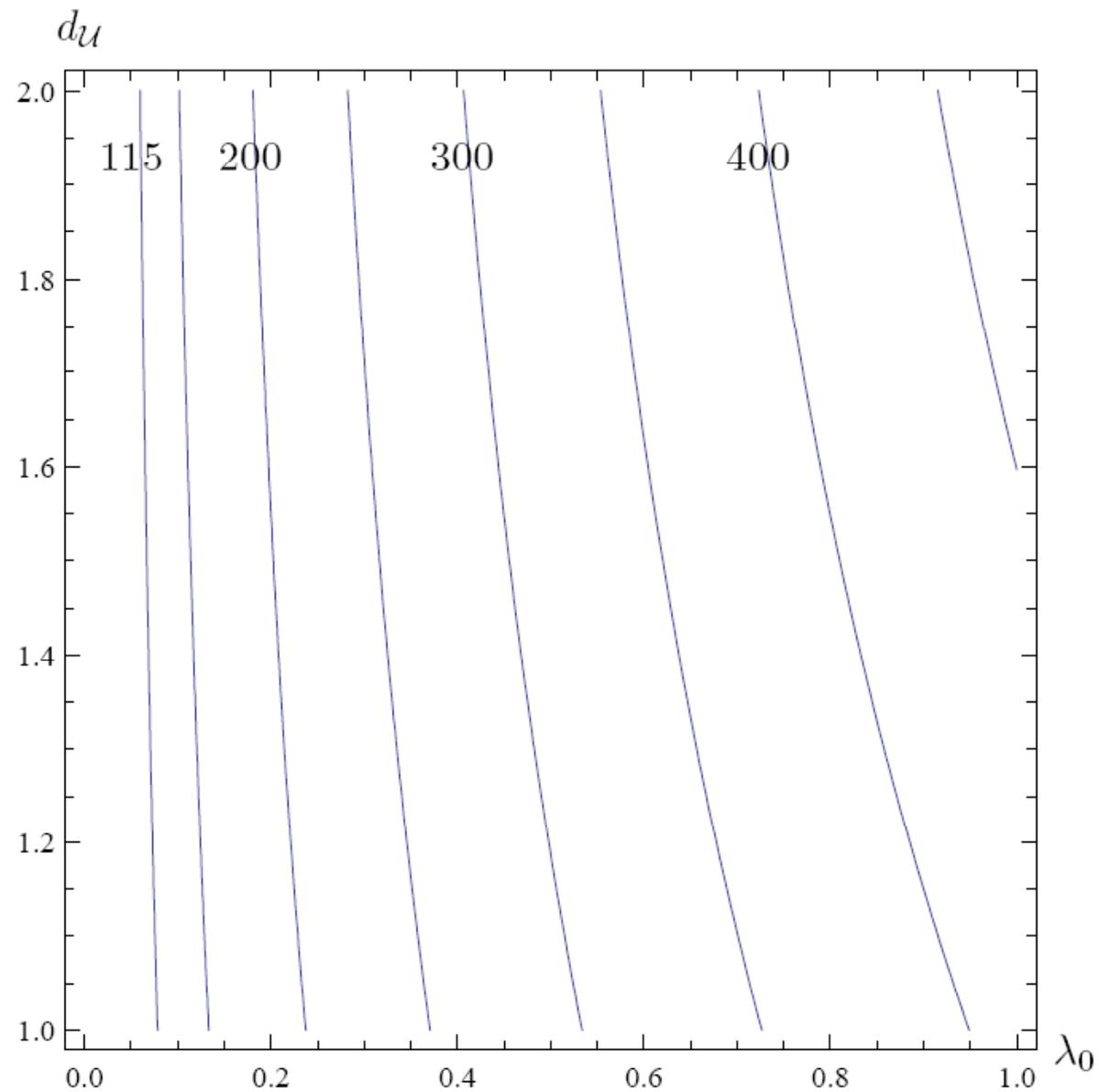
$$\frac{m_h^2}{\rho_0^2} = 2\sqrt{\lambda_0\lambda_1} \left[ 1 + \frac{\epsilon}{2} \left( \frac{\sqrt{\lambda_0} - \sqrt{\lambda_1}}{\sqrt{\lambda_0} + \sqrt{\lambda_1}} \right)^2 \right]$$

$$\frac{m_\ell^2}{\rho_0^2} = 2\sqrt{\lambda_0\lambda_1} \left[ \frac{2\epsilon\sqrt{\lambda_0\lambda_1}}{(\sqrt{\lambda_0} + \sqrt{\lambda_1})^2} \right] .$$

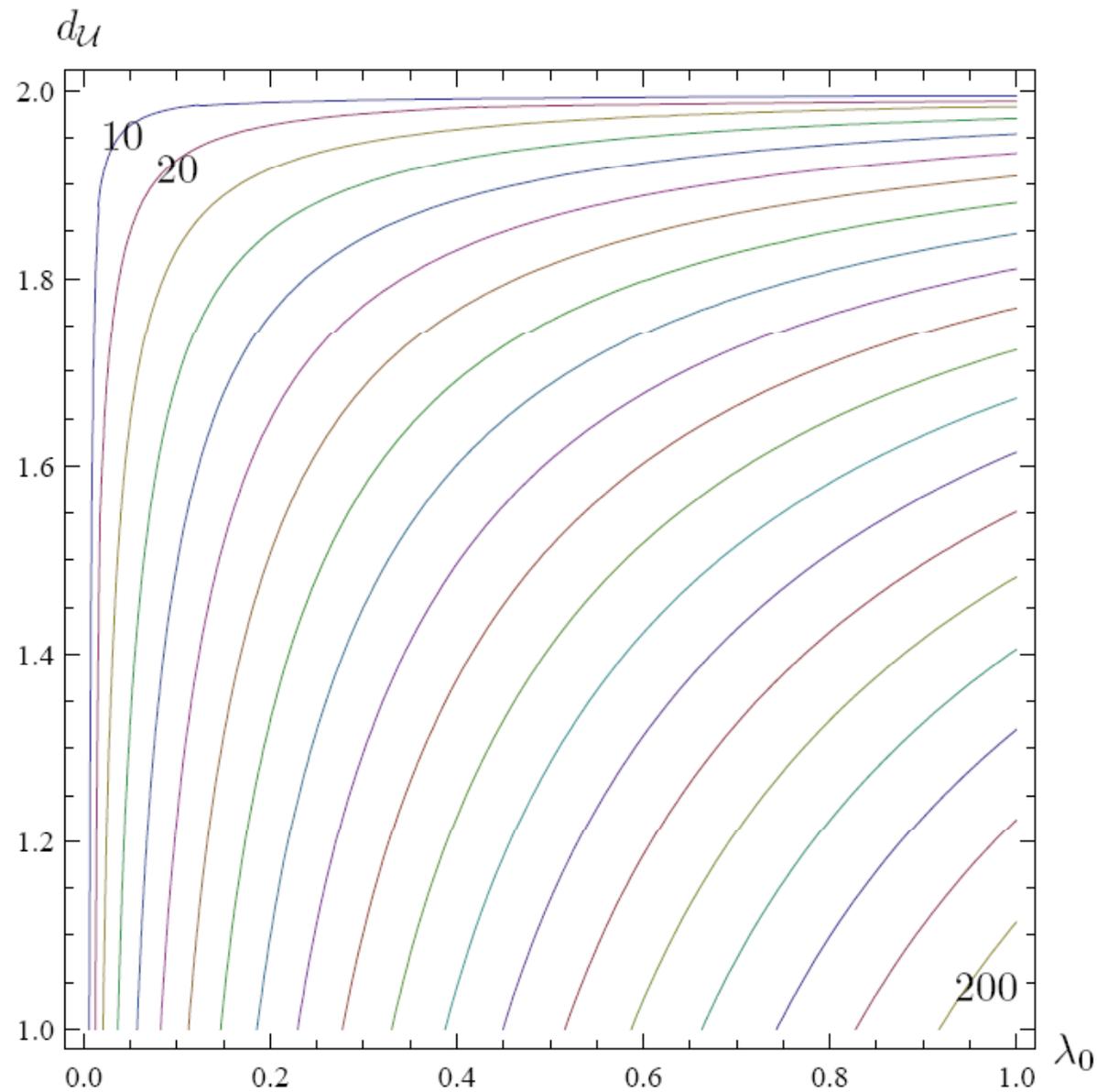
$m_h$  : Higgs boson

$m_\ell$  : New light scalar

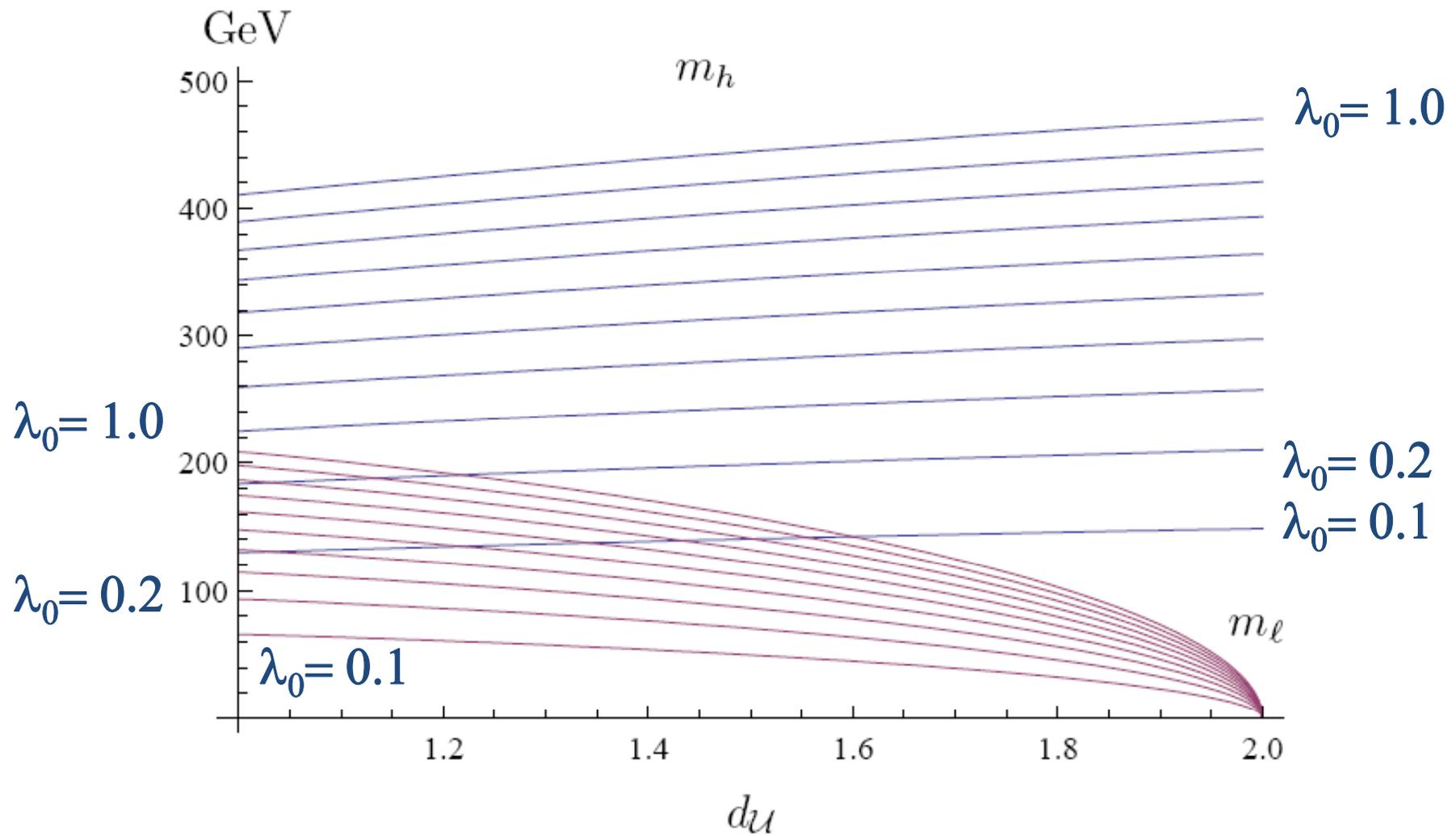
# Contours for $m_h$ (GeV)



# Contours for $m_\ell$ (GeV)



# $m_{h,\ell}$ vs $d\mu$



# Conclusions

- Unparticles inspire a new scalar potential.
- EWSB occurs when the unparticle sector begins to interact with the SM sector.
- If the hidden sector were not scale inv., the EWSB occurs via CW mechanism.
- When  $d_U$  starts to depart from 2, a new scale is introduced to break EWS.
- After the EWSB, a new light scalar  $\lesssim 210\text{GeV}$  appears.

**Back Up Slides**

